

Recovery of differentiation/integration compatibility of meshless operators via local adaptation of the point cloud in the context of nodal integration

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Virtual Product Performance Management





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- 2 Re-interpretation of "background pressure"-like corrections
- 3 Scaling of the parameter of the method
 - Conclusion and future work



Cloud of points : CBondary nodes $\partial C \subset C$





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Meshless operators :

• Nodal volume quadrature : $\oint_{\mathcal{C}} f = \sum_{i \in \mathcal{C}} V_i f_i$

• Boundary quadrature :
$$\oint_{\partial \mathcal{C}} f = \sum_{i \in \partial \mathcal{C}} f_i \Gamma_i$$



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• Meshless gradient : $V_i \nabla_i f = \sum_{j \in \mathcal{N}(i)} \mathbf{A}_{i,j} f_j$



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• Meshless gradient : $V_i \nabla_i f = \sum_{j \in \mathcal{N}(i)} \mathbf{A}_{i,j} f_j \begin{cases} \nabla_i^{\text{SPH}} f = \sum_j V_j \nabla W(\mathbf{x}_j - \mathbf{x}_i) f_j \\ \nabla_i^{\text{R0}} f = \sum_j V_j \nabla W(\mathbf{x}_j - \mathbf{x}_i) (f_j - f_i) \\ \nabla_i^{\text{R1}} f = \sum_j V_j B_i \nabla W(\mathbf{x}_j - \mathbf{x}_i) (f_j - f_i) \end{cases}$

$$\nabla_i^{\text{MLS}} f = MI \text{ S interpolation at } \mathbf{x}_i$$



Integration by parts formula

$$\int_{\Omega} f \nabla \cdot \mathbf{v} + \mathbf{v} \cdot \nabla f \, \mathrm{d}V = \int_{\partial \Omega} f \mathbf{v} \cdot \mathrm{d}\mathbf{S}$$

Discrete counterpart : definition of a dual gradient ∇^*

$$\oint_{\mathcal{C}} f \nabla \cdot \mathbf{v} + \mathbf{v} \cdot \nabla^* f = \oint_{\partial \mathcal{C}} f \mathbf{v}$$



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Discrete counterpart : definition of a dual gradient \mathbb{V}^*

$$\oint_{\mathcal{C}} f \nabla \cdot \mathbf{v} + \mathbf{v} \cdot \nabla^* f = \oint_{\partial \mathcal{C}} f \mathbf{v}$$

Explicit formula for the dual gradient

$$V_i \mathbb{V}_i^* f = \sum_{j \in \mathcal{N}(i)} (-\mathbf{A}_{j,i} + \delta_{i,j} \mathbf{\Gamma}_i) f_j$$

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Continuous weak formulation

Find $u \in \mathcal{H}^1(\Omega)$ such that :

$$\begin{cases} \int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} sv \quad \forall v \in \mathcal{H}_0^1(\Omega) \\ u_{|\partial\Omega} = u_0 \end{cases}$$

Discrete weak formulation

Find $u: \mathcal{C} \to \mathbb{R}$ such that :

$$\begin{cases} \oint_{\mathcal{C}} \nabla u \cdot \nabla v = \oint_{\mathcal{C}} sv \quad \forall v : \mathcal{C} \setminus \partial \mathcal{C} \to \mathbb{R} \\ u_{|\partial \mathcal{C}} = u_0 \end{cases}$$

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Find $u: \mathcal{C} \to \mathbb{R}$ such that :

$$\begin{cases} -\mathbb{V}_i^* \cdot \mathbb{\nabla} u = s_i & \forall i \in \mathcal{C} \backslash \partial \mathcal{C} \\ u_{|\partial \mathcal{C}} = u_0 \end{cases}$$



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Necessary conditions for the linear patch test

•
$$\nabla x = \mathbf{I}_d$$

•
$$\nabla^* 1 = 0$$

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Necessary conditions for the linear patch test

•
$$\nabla x = \mathbf{I}_{a}$$

• $\nabla^* 1 = 0$ Compatibility !



Find $u: \mathcal{C} \to \mathbb{R}$ such that :

$$\begin{cases} -\mathbb{V}_i^* \cdot \mathbb{\nabla} u = s_i & \forall i \in \mathcal{C} \backslash \partial \mathcal{C} \\ u_{|\partial \mathcal{C}} = u_0 \end{cases}$$

Necessary conditions for the linear patch test

•
$$\nabla x = \mathbf{I}_d$$

•
$$\nabla^* 1 = 0$$

 \Leftrightarrow Discrete Stokes formula :

$$\oint_{\mathcal{C}} \nabla u = \oint_{\partial \mathcal{C}} u$$

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Corrected gradient

$$\mathbb{V}_i^c u = \mathbb{V}_i u + \sum_{j \in \mathcal{N}(i)} \boldsymbol{\mu}_{i,j} (u_j - u_i - \mathbb{V}_i u \cdot (\mathbf{x}_j - \mathbf{x}_i))$$

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Corrected gradient

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Correction preserves linear consistency :

$$\forall \boldsymbol{\mu}_{i,j}, \ \nabla \mathbf{x} = \mathbf{I}_d \Rightarrow \nabla^c \mathbf{x} = \mathbf{I}_d$$

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Correction equations

Solve $\nabla^{c*1} = 0$ for $\mu_{i,i}$ given ∇

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A diffusion test case





Halton sequences

$$s = 20\pi^2 \sin(2\pi x) \sin(4\pi y)$$

$$u = \sin(2\pi x) \sin(4\pi y)$$





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🙂 Cons

- An additional global linear system to solve
 ... every timestep !
- No explicit formula for the gradient

Approximate compatibility









Observation : Keep $||\nabla^*1|| = O(1)$ instead of the usual $||\nabla^*1|| = O(h^{-1})$ \Rightarrow recover almost second order convergence !

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Another idea :



Go from :



$$\left\|\nabla^{R1*1}\right\| \approx 36.3$$

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Another idea :





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h-convergence for the diffusion equation





$$\mathbb{V}_i^{\text{R0}} f = -\sum_{j \in \mathcal{C}} V_j \nabla W_h(\mathbf{x}_j - \mathbf{x}_i) (f_j - f_i)$$

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h-convergence for the diffusion equation





$$\nabla_i^{\text{R0}} f = -\sum_{j \in \mathcal{C}} V_j \nabla W_h(\mathbf{x}_j - \mathbf{x}_i) (f_j - f_i)$$
$$\nabla_i^{\text{R0}} \mathbf{x} \neq \mathbf{I}_d$$

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h-convergence for the diffusion equation





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Statics \Rightarrow A gradient-specific "remeshing" procedure



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 $Dynamics \Rightarrow A \text{ gradient-specific ALE source term}$



Statics \Rightarrow A gradient-specific "remeshing" procedure

Dynamics \Rightarrow A gradient-specific ALE source term

The idea is not new !

- XSPH (Monaghan)
- Background pressure
- Transport-velocity formulation
- Fickian-based shifting

[Adami 2013] [Lind 2011]



Statics \Rightarrow A gradient-specific "remeshing" procedure

Dynamics \Rightarrow A gradient-specific ALE source term

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[Adami 2013] [Lind 2011]



$$m_i \frac{\mathrm{d} \mathbf{v}_i}{\mathrm{d} t} = \mathbf{f}_i$$
$$\frac{\mathrm{d} \mathbf{x}_i}{\mathrm{d} t} = \mathbf{v}_i$$

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$$m_i \frac{\mathrm{d}\mathbf{v}_i}{\mathrm{d}t} = \mathbf{f}_i \quad -V_i \nabla_i^* P_u$$
$$\frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}_i$$

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$$m_i \frac{\mathrm{d}\mathbf{v}_i}{\mathrm{d}t} = \mathbf{f}_i \quad -V_i P_u \nabla_i^* \mathbf{1}$$
$$\frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}_i$$

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$$m_i \frac{\mathrm{d} \mathbf{v}_i}{\mathrm{d} t} = \mathbf{f}_i$$
$$\frac{\mathrm{d} \mathbf{x}_i}{\mathrm{d} t} = \mathbf{v}_i \quad -V_i \alpha P_u \mathbb{V}_i^* \mathbf{1}$$

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Classical and renormalized SPH of order 0

$$V_i \mathbb{V}_i^{SPH*} 1 = \frac{1}{2} V_i \mathbb{V}_i^{R0*} 1 = V_i \mathbb{V}_i^{SPH} 1 = -\sum_{j \in \mathcal{N}(i)} V_i V_j \nabla W(\mathbf{x}_j - \mathbf{x}_i)$$

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Constant quadrature : V_i independent of $(\mathbf{x}_j)_{j \in C}$

$$V_i \mathbb{V}_i^{SPH*} 1 = \frac{\partial}{\partial \mathbf{x}_i} V_{tot}$$

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Classical and renormalized SPH of order 0

$$V_i \mathbb{V}_i^{SPH*} 1 = \frac{1}{2} V_i \mathbb{V}_i^{R0*} 1 = V_i \mathbb{V}_i^{SPH} 1 = -\sum_{j \in \mathcal{N}(i)} V_i V_j \nabla W(\mathbf{x}_j - \mathbf{x}_i)$$

Constant quadrature : V_i independent of $(\mathbf{x}_j)_{j \in C}$

$$\begin{split} V_i \nabla_i^{SPH*} 1 &= \frac{\partial}{\partial \mathbf{x}_i} V_{tot} \\ V_{tot} &= \sum_k V_k \sum_j V_j W(\mathbf{x}_j - \mathbf{x}_k) \\ &= \oint_{\mathcal{C}} < 1 > \end{split}$$

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Density-based nodal quadrature

$$V_{i} = \frac{1}{\sum_{j} W(\mathbf{x}_{j} - \mathbf{x}_{i})}$$
$$V_{i} \mathbb{V}_{i}^{Adami*} 1 = -\sum_{j \in \mathcal{N}(i)} (V_{i}^{2} + V_{j}^{2}) \nabla W(\mathbf{x}_{j} - \mathbf{x}_{i}) \quad \text{[Adami 2013]}$$

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Density-based nodal quadrature

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$$V_{i} \nabla_{i}^{Adami*} 1 = -\sum_{j \in \mathcal{N}(i)} (V_{i}^{2} + V_{j}^{2}) \nabla W(\mathbf{x}_{j} - \mathbf{x}_{i}) \quad \text{[Adami 2013]}$$

$$= \frac{\partial}{\partial \mathbf{x}_{i}} V_{tot}$$

$$V_{tot} = \sum_{k} V_{k} = \oint_{\mathcal{C}} 1$$

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Background pressure and compatibility



ALE source term

$$\frac{\mathrm{d}\mathbf{x}_{i}}{\mathrm{d}t} = -\gamma V_{i} \nabla^{*} 1$$

$$\mathcal{L}((\mathbf{x}_{i})_{i \in \mathcal{C}}) = -V_{tot}$$

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Background pressure and compatibility



ALE source term

$$\frac{\mathrm{d}\mathbf{x}_{i}}{\mathrm{d}t} = -\gamma V_{i} \nabla^{*} 1$$

$$\mathcal{L}((\mathbf{x}_{i})_{i \in \mathcal{C}}) = -V_{tot}$$

 \Rightarrow Convergence towards a node distribution where $\frac{\partial \mathcal{L}}{\partial \mathbf{x}_i} = V_i \nabla_i^* 1 = 0$

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Background pressure and compatibility





Background pressure

$$\begin{cases} \frac{\mathrm{d}\mathbf{v}_i}{\mathrm{d}t} = -\alpha \mathbf{v}_i - \gamma V_i \nabla^* 1\\ \frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}_i \end{cases}$$
$$(\mathbf{x}_i, \mathbf{v}_i)_{i \in \mathcal{C}}) = \frac{1}{2} \sum \|\mathbf{v}_i\|^2 - \gamma V_{tot}$$

 \Rightarrow Convergence towards a node distribution where $\frac{\partial \mathcal{L}}{\partial \mathbf{x}_i} = V_i \nabla_i^* 1 = 0$

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|------------------|----------|
|------------------|----------|

The simplest dynamic problem : advection



$$\frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}(\mathbf{x}_i) \quad \forall i \in \mathcal{C} \backslash \partial \mathcal{C}$$

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The simplest dynamic problem : advection







Taylor-Green vortices

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The simplest dynamic problem : advection







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Approximate advection



$$\begin{cases} \frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}(\mathbf{x}_i) - \mathbf{c}_i \quad \forall i \in \mathcal{C} \backslash \partial \mathcal{C} \\ \mathbf{c}_i = \gamma \nabla^{*\mathrm{SPH}} 1 \end{cases}$$

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Approximate advection



$$\begin{cases} \frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}(\mathbf{x}_i) - \mathbf{c}_i & \forall i \in \mathcal{C} \backslash \partial \mathcal{C} \\ \mathbf{c}_i = \gamma \nabla^{*\mathrm{SPH}} \mathbf{1} \end{cases}$$

The ALE speed

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Approximate advection



$$\begin{cases} \frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}(\mathbf{x}_i) - \mathbf{c}_i \quad \forall i \in \mathcal{C} \backslash \partial \mathcal{C} \\ \mathbf{c}_i = \gamma \nabla^{*\mathrm{SPH}} \mathbf{1} \end{cases}$$



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Influence of h





Influence of h





$$\begin{aligned} \|\nabla^{SPH*}1\| \propto h^{0.35} \\ \|\mathbf{c}\| \propto h^{0.35} \end{aligned}$$





$$h = 3.3 \cdot 10^{-2}$$

$$h = 2 \cdot 10^{-2}$$

$$h = 1.4 \cdot 10^{-2}$$

$$h = 1 \cdot 10^{-2}$$

$$h = 6.6 \cdot 10^{-3}$$

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$$\begin{aligned} \|\nabla^{SPH*1}\| \propto h^{0.35} \gamma^{-0.65} \\ \|\mathbf{c}\| \propto h^{0.35} \gamma^{0.35} \end{aligned}$$

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Scaling of parameters



$$\frac{\|\nabla^{SPH*}1\| \propto h^{0.35} \gamma^{-0.65}}{\|\mathbf{c}\| \propto h^{0.35} \gamma^{0.35}}$$

$$\gamma = \mathcal{O}(\sqrt{h}) \Rightarrow \begin{cases} \| \mathbb{V}^{SPH*1} \| = \mathcal{O}(1) \\ \| \mathbf{c} \| = \mathcal{O}(\sqrt{h}) \end{cases}$$

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Fick-based diffusion [Lind 2011]



$$\begin{cases} \frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}(\mathbf{x}_i) - \mathbf{c}_i \quad \forall i \in \mathcal{C} \setminus \partial \mathcal{C} \\ \mathbf{c}_i = \gamma \nabla^{*\mathrm{SPH}} 1 \end{cases}$$

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... then uses \mathbb{V}^{R1} for Poisson equation !

Fick-based diffusion [Lind 2011]

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$$\begin{cases} \frac{\mathrm{d}\mathbf{x}_i}{\mathrm{d}t} = \mathbf{v}(\mathbf{x}_i) - \mathbf{c}_i \quad \forall i \in \mathcal{C} \backslash \partial \mathcal{C} \\ \mathbf{c}_i = \gamma \nabla^{*\mathrm{SPH}} \mathbf{1} \end{cases}$$



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Summary

• Re-interpretation of several background-pressure like procedures within an operator framework

- Proved convergence of the method to compatible steady state for several existing gradient
- We proposed a proper scaling for the parameter of the method



Summary

- Re-interpretation of several background-pressure like procedures within an operator framework
- Proved convergence of the method to compatible steady state for several existing gradient
- We proposed a proper scaling for the parameter of the method

Ongoing and future work

- Clarify the role of volume related functionals
- Extend methodology to handle other boundary treatments





Thanks for your attention !

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